

cules treated in the one configuration approximation. Likewise continuum orbitals could be invoked in situations where the optimum forms of the molecular orbitals of a molecular fragment change substantially during the process of molecule formation.

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The Crystal Structure of $\text{Yb}_{11}\text{Sb}_{10}$

Sir:

The compound previously designated as $\alpha\text{-Yb}_5\text{Sb}_4$ is the low-temperature form of one of the intermediate phases occurring in the Yb-Sb system investigated by Bodnar and Steinfink.¹ The stoichiometry of the compound had not previously been established with certainty, although a single phase was observed on an X-ray powder diffraction pattern when the elements reacted in a 5:4 ratio. The lattice parameters and space group of this compound indicated that it was isostructural with one of the monogermanide phases reported by Tharp, Smith, and Johnson.² A later single-crystal analysis by Smith, Johnson, and Tharp³ established the structure of this new rare earth germanide phase and showed that the composition corresponded to an 11:10 ratio of the elements.

Procedure.—The preparation of this compound has been described previously.¹ The material is spongy,

Atom 1	Atom 2	Distance, Å	Esd, Å
Yb(1)	Sb(4)	3.125	0.003
	Sb(2)	3.159	0.006
	Sb(5)	3.413	0.007
	Sb(1)	3.420	0.005
	Sb(5)	3.527	0.007
	Yb(4)	3.701	0.004
	Yb(4)	3.819	0.004
	Yb(3)	3.920	0.005
	Yb(2)	3.910	0.009
	Yb(1)	4.223	0.004
Yb(2)	Sb(1)	2.933	0.010
	Sb(3)	3.464	0.018
	Sb(2)	3.556	0.010
	Sb(5)	3.672	0.011
	Yb(4)	4.319	0.009
Yb(3)	Yb(2)	3.963	0.013
	Sb(3)	3.491	0.016
	Sb(5)	3.465	0.008
Yb(4)	Yb(4)	4.119	0.005
	Sb(5)	3.148	0.007
	Sb(1)	3.165	0.005
	Sb(3)	3.393	0.016
	Sb(4)	3.176	0.003
Sb(1)	Yb(4)	3.881	0.004
	Sb(1)	3.572	0.006
	Sb(3)	2.998	0.022
	Sb(5)	3.280	0.017
Sb(5)	Sb(5)	2.936	0.009

^a Esd's include only uncertainties in positional parameters.

density was redetermined by immersing the specimen in benzene and evacuating the system so that boiling occurred and benzene vapors could penetrate into the pores to displace the air. When the system was opened

TABLE I
FINAL ATOMIC PARAMETERS ($\times 10^4$) FOR $\text{Yb}_{11}\text{Sb}_{10}$ ^a

Atom	Wyckoff notation	x	y	z	B_{11}	B_{22}	B_{33}	B_{12}	B_{13}	B_{23}
Yb(1)	16n	0	2518 (2)	1886 (1)	38 (1)	22 (1)	6 (0)	0	0	-1 (1)
Yb(2)	8h	1671 (5)	1671 (5)	0	71 (4)	71 (4)	14 (1)	40 (3)	0	0
Yb(3)	4e	0	0	3374 (2)	32 (4)	32 (4)	11 (1)	0	0	0
Yb(4)	16n	0	3364 (2)	3974 (1)	26 (1)	36 (1)	4 (0)	0	0	-1 (1)
Sb(1)	8i	3494 (4)	0	0	26 (3)	45 (3)	5 (1)	0	0	0
Sb(2)	4e	0	0	1283 (3)	23 (5)	23 (5)	8 (2)	0	0	0
Sb(3)	8h	3736 (9)	3736 (9)	0	104 (8)	104 (8)	19 (2)	-1 (5)	0	0
Sb(4)	4d	0	1/2	1/4	25 (4)	25 (4)	3 (1)	0	0	0
Sb(5)	16m	2062 (4)	2062 (4)	3245 (2)	35 (2)	35 (2)	8 (1)	6 (2)	-1 (2)	-1 (2)

^a The esd's $\times 10^4$ are given in parentheses. The form of the thermal ellipsoids is $\exp[-(B_{11}h^2 + B_{22}k^2 + B_{33}l^2 + 2B_{12}hk + 2B_{13}hl + 2B_{23}kl)]$.

and single crystals were difficult to find. The single crystal used in this investigation was an irregular fragment with encompassing dimensions of about $0.06 \times 0.07 \times 0.10$ mm and with a shape which could be approximated by six planar faces.

The calculated density assuming 4 units of $\text{Yb}_{11}\text{Sb}_{10}$ per unit cell is 8.6 g/cm^3 . Density measurements on the spongy material by loss of weight in acetone resulted in a suspiciously low value of 7.6 g/cm^3 , probably because acetone failed to displace entrapped air. The

to the atmosphere, the vapors condensed to liquid in the pores and additional liquid was sucked in so that a true weight loss was determined. Several such measurements yielded values ranging from 8.11 to 8.24 g/cm^3 , with an average value of $8.17 (7) \text{ g/cm}^3$, considerably closer to the calculated value for the stoichiometry $\text{Yb}_{11}\text{Sb}_{10}$.

Weissenberg and precession photographs, made with Cu $K\alpha$ and Mo $K\alpha$ radiation, respectively, showed that the diffraction symmetry was $4/mmm$, that the crystal was mounted along an axis which was labeled b , and that hkl diffraction spectra were present only when $h + k + l = 2n$ with no other systematic absences being

(1) R. E. Bodnar and H. Steinfink, *Inorg. Chem.*, **6**, 327 (1967).

(2) A. G. Tharp, G. S. Smith, and Q. Johnson, *Acta Crystallogr.*, **20**, 583 (1966).

(3) G. S. Smith, Q. Johnson, and A. G. Tharp, *ibid.*, **23**, 640 (1967).

equivalents for an independent reflection were statistically acceptable, only the acceptable ones were used in calculating the average quantities. In the other cases all equivalents were used. The averaged reflections were combined with single reflections for which no equivalents were available in the data set to obtain a final set of 1370 independent reflections, of which 997 were statistically nonzero.

Determination of the Structure.—In the refinements discussed below, R is the conventional discrepancy coefficient, $R = \Sigma |F_o| - |F_c| / \Sigma |F_o|$. The function minimized was $\Sigma [w(|F_o| - |F_c|)^2]$, where w is $[\sigma(F_o)]^{-2}$. The computer program used in the refinements was XFLS, a modified version of ORFLS.⁶ Dispersion-corrected scattering factors incorporating the real and imaginary terms for Yb and Sb as given in ref 7 for Mo $K\alpha$ were used throughout.

Beginning with a trial model consisting of the $\text{Ho}_{11}\text{Ge}_{10}$ parameters, eight cycles of full-matrix least-squares refinement of the scale factor, positional parameters, and isotropic thermal parameters resulted in an R value of 0.13. The thermal parameters after this refinement all fell within the range of 0.3–1.0 \AA^2 with the exception of Sb(3), for which the thermal parameter was 3.9 \AA^2 . Five cycles of anisotropic refinement then resulted in $R = 0.080$, with the thermal parameters for Sb(3) remaining quite high relative to those for the other atoms in the structure.

The high thermal parameter for Sb(3) and the large value of the calculated density for $\text{Yb}_{11}\text{Sb}_{10}$ might be an indication that the Sb(3) site is partially occupied. A stoichiometry of $\text{Yb}_{11}\text{Sb}_9$ is obtained on the assumption that the occupant is $1/2\text{Sb}$. The calculated density for

this formula is 8.3 g/cm^3 . A full-matrix refinement on this basis reduced the temperature factors of Sb(3) to values comparable with those of the other atoms, as expected, but R increased to 0.084. The formula $\text{Yb}_{11}\text{Sb}_{10}$ was retained for all further calculations but the possibility that a deviation from the 11:10 composition exists cannot be ruled out.

The final parameters for $\text{Yb}_{11}\text{Sb}_{10}$ are shown in Table I, the interatomic distances are shown in Table II, and the final observed amplitudes and calculated structure factors are presented in Table III. The standard deviation of an observation of unit weight is 2.67. A difference Fourier synthesis was calculated and contained no peaks representing more than about 0.4 $e^{-}/\text{\AA}^3$.

Discussion of the Structure.—The spatial parameters of " $\alpha\text{-Yb}_5\text{Sb}_4$ " show that it is isostructural with $\text{Ho}_{11}\text{Ge}_{10}$. A detailed discussion of this structure has been given by Smith, *et al.*³ All of the bond distances in Table II fall within an acceptable range for the atoms involved. The same type of bonding is expected in this compound as in $\text{Ho}_{11}\text{Ge}_{10}$, because valence numbers of 3 and 4 for Ho and Ge, respectively, would provide the same number of electrons for bonding as valence numbers of 2 and 5 for Yb and Sb, respectively. Magnetic susceptibility data indicate that nearly all of the Yb is divalent in this phase.⁸

(8) R. E. Bodnar, H. Steinfink, and K. S. V. L. Narasimhan, *J. Appl. Phys.*, **39**, 1485 (1968).

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(6) W. R. Busing, K. O. Martin, and H. A. Levy, USAEC Report ORNL-TM 305, Oak Ridge National Laboratory, Oak Ridge, Tenn., 1962.

(7) "International Tables for X-Ray Crystallography," Vol. III, Kynoch Press, Birmingham, England, 1968.

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